

On infinite energy solutions of Schrödinger-type Equations with a nonlocal term

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Joint work with Ademir Pastor (UNICAMP)

October 30 - November 1, 2013

First Workshop on Nonlinear Dispersive Equations

Schrödinger-type Equations with a nonlocal term

Initial value problem (IVP) associated with
Schrödinger-type equations of the form

$$\begin{cases} i\partial_t u + Lu = \chi|u|^\rho u + bE(|u|^\rho)u, \\ u(x, 0) = u_0(x). \end{cases} \quad (1)$$

- $(x, t) \in \mathbb{R}^n \times \mathbb{R}$, $n \geq 1$,
- $u = u(x, t)$ is a complex-valued function,
- χ and b are real constants, ρ is a positive real number
- L and E are linear operators.

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Some examples

- Schrödinger equation

$$i\partial_t u + \Delta u = \chi|u|^\rho u. \quad (2)$$

- Davey-Stewartson system ($n \geq 2$, $m > 0$)

$$\begin{cases} i\partial_t u + \delta\partial_{x_1}^2 u + \sum_{j=2}^n \partial_{x_j}^2 u = \chi|u|^\rho u + bu\partial_{x_1}\varphi, \\ \partial_{x_1}^2 \varphi + m\partial_{x_2}^2 \varphi + \sum_{j=3}^n \partial_{x_j}^2 \varphi = \partial_{x_1}(|u|^\rho), \\ u(x, 0) = u_0(x). \end{cases} \quad (3)$$

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- To $m > 0$, system (3) becomes

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where

$$\widehat{E(f)}(\xi) = \frac{\xi_1^2}{\xi_1^2 + m\xi_2^2 + \sum_{j=3}^n \xi_j^2} \hat{f}(\xi) = p(\xi)\hat{f}(\xi).$$

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In [B]

$$i\partial_t u + \delta \partial_{x_1}^2 u + \sum_{j=2}^n \partial_{x_j}^2 u = \chi |u|^\rho u + buE(|u|^\rho)$$

- Lorentz spaces:

$$L^{p,\infty}(\mathbb{R}^n) = \{f; \|f\|_{L^{p,\infty}(\mathbb{R}^n)} := \sup_{\lambda > 0} \lambda \alpha(\lambda, f)^{1/p} < \infty\}$$

$$\alpha(\lambda, f) = \mu(\{x \in \mathbb{R}^n; |f(x)| > \lambda\}),$$

- Global in time solutions (\Rightarrow self-similar solutions)

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Pablo Silva, Lucas Ferreira, Elder Roa ([SFR]):

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- $E_\alpha = \{u; \|u\|_\alpha = \sup_{-\infty < t < +\infty} |t|^{\alpha/2} \|u(t)\|_{L(\rho+2, \infty)} < \infty\}$,
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Davey-Stewartson ($m > 0$)

$$i\partial_t u + \delta\partial_{x_1}^2 u + \sum_{j=2}^n \partial_{x_j}^2 u = \chi|u|^\rho u + buE(|u|^\rho)$$

Extends the results of [SFR] to the D-S system.

$$L = \delta\partial_{x_1}^2 + \sum_{j=2}^n \partial_{x_j}^2 \rightarrow L$$

$$\widehat{E(f)}(\xi) = \frac{\xi_1^2}{\xi_1^2 + m\xi_2^2 + \sum_{j=3}^n \xi_j^2} \hat{f}(\xi) = p(\xi)\hat{f}(\xi) \rightarrow E$$

To give sufficient conditions on the operators L and E that allow to extend the results to the IVP (1).

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To give sufficient conditions on the operators L and E that allow to extend the results to the IVP (1) .

Conditions on L and E

L is a pseudo-differential operator defined via its Fourier transform by

$$\widehat{Lu}(\xi) = q(\xi)\widehat{u}(\xi), \quad (4)$$

(H1) the function q is real and homogeneous of degree d , that is,

$$q(\lambda\xi) = \lambda^d q(\xi), \quad \lambda > 0.$$

(H2) The function $G(x) = \int_{\mathbb{R}^n} e^{i(x\xi+q(\xi))} d\xi$ belongs to $L^\infty(\mathbb{R}^n)$.

(H3) E is bounded from $L^{(p,\infty)}(\mathbb{R}^n)$ to itself, for all p satisfying $1 < p < \infty$.

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- Self-similar solutions + conditions,
- Our main results,
- Applications.

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Self-similar solutions

If $u(x, t)$ is a solution of (1) so is $u_\lambda(x, t) = \lambda^{d/\rho}u(\lambda x, \lambda^d t)$, for any $\lambda > 0$.

important!(H1)- $q(\lambda\xi) = \lambda^d q(\xi)$

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Supposing

- $\exists!$ of solutions to the IVP problem (1)
- u a self-similar solution

We must have

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i.e., $u_0(x) = \lambda^{d/\rho} u_0(\lambda x)$ (u_0 is homogeneous).

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Main Results

Integral equivalent formulation to the IVP (1)

$$u(t) = U(t)u_0 + i \int_0^t U(t-s)(\chi|u|^\rho u + buE(|u|^\rho))(s)ds, \quad (5)$$

where $U(t)u_0$ is the solution of the linear problem

$$\begin{cases} i\partial_t u + Lu = 0, \\ u(x, 0) = u_0(x) \end{cases} \quad (x, t) \in \mathbb{R}^n \times \mathbb{R}, \quad (6)$$

that is,

$$U(t)u_0(x) = \int_{\mathbb{R}^n} e^{i(x\xi + tq(\xi))} \widehat{u}_0(\xi) d\xi. \quad (7)$$

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Theorem (Global Existence)

- $1 < \rho < \infty$ and $\frac{\rho+2}{\rho+1} < \frac{n\rho}{d} < \rho + 2$,
- ϕ is a distribution satisfying $\|U(t)\phi\|_\alpha \leq \epsilon$, where $0 < \epsilon \ll 1$.

Then

- The integral equation (5) has a unique solution $u \in E_\alpha$ satisfying $\|u\|_\alpha \leq 2\epsilon$,
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Lemma (A)

- $1 < p < 2$,
- p' such that $\frac{1}{p} + \frac{1}{p'} = 1$.

Then there exists a constant $C = C(n, p) > 0$ such that

- $\|U(t)\phi\|_{L(p', \infty)} \leq C|t|^{-\frac{n}{d}(\frac{2}{p}-1)} \|\phi\|_{L(p, \infty)}$,
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$$B(u) = i \int_0^t U(t-s)(\chi|u|^\rho u + buE(|u|^\rho))(s)ds.$$

Lemma (B)

- $1 < \rho < \infty$ and $\frac{\rho+2}{\rho+1} < \frac{n\rho}{d} < \rho + 2$

Then there exists a positive constant K_α such that

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- Lemma (A) + **(H3)** E is bounded from $L^{(p,\infty)}(\mathbb{R}^n)$ to itself, for all p satisfying $1 < p < \infty$.

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$$B(u) = i \int_0^t U(t-s)(\chi|u|^\rho u + buE(|u|^\rho))(s)ds.$$

Lemma (B)

- $1 < \rho < \infty$ and $\frac{\rho+2}{\rho+1} < \frac{n\rho}{d} < \rho + 2$

Then there exists a positive constant K_α such that

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Lemma (B) + hypothesis $\|U(t)\phi\|_\alpha \leq \epsilon$ implies

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We need to ask ϕ homogeneous

- $\phi(\lambda x) = \lambda^{-\frac{d}{\rho}} \phi(x), \forall \lambda > 0.$
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- the solution u obtained in Theorem Global Existence is self-similar.

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Other Results-Scattering

Theorem (Scattering)

$$\|u(t) - U(t)u_{\pm}\|_{L(\rho+2,\infty)} \leq C|t|^{-\frac{\alpha}{d}} \|u\|_{\alpha}^{\rho+1}, \quad t \neq 0.$$

$$u \longrightarrow v(t) = U(t)\phi + B(v)$$

$$U(t)u_{\pm} \longrightarrow \begin{cases} i\partial_t v + Lv = 0, \\ v(0) = u_{\pm}, \end{cases}$$

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Standard NLS Equation ($n \geq 1$)

$$i\partial_t u + \Delta u = \chi u|u|^\rho, \quad (x, t) \in \mathbb{R}^n \times \mathbb{R}, \quad n \geq 1,$$

[SFR]

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Nonelliptic NLS Equation ($n \geq 1$)

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(H1) the function $q(x) = x_1^2 - x_2^2$ homogeneous of degree 2,

(H2) $G(x) = \int_{\mathbb{R}^n} e^{i(x\xi + q(\xi))} d\xi \in L^\infty(\mathbb{R}^n)$ see [GS1].

b=0 (**(H3)** E is bounded from $L^{(p,\infty)}(\mathbb{R}^n)$ to itself, for all p satisfying $1 < p < \infty$.)

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$$\begin{cases} i\partial_t u + \frac{\omega_{kk}}{2}\partial_x^2 u + \frac{\omega_{\ell\ell}}{2}\partial_y^2 u + \frac{\omega_{nn}}{2}\partial_z^2 u + \omega_{nk}\partial_{xz}^2 u = -uQ, \\ \partial_x^2 Q + \partial_y^2 Q = \nu\partial_y^2 |u|^\rho, \end{cases}$$

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