# Almost sure global well-posedness for the cubic wave equation 

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## Wave equation

We consider the cubic wave equation on $\mathbb{R}^{3}$ :

$$
\left\{\begin{array}{l}
\partial_{t}^{2} f-\Delta f+f^{3}=0 \\
f_{\mid t=0}=f_{0}, \partial_{t} f_{t=0}=f_{1}
\end{array}\right.
$$

The critical exponent of this equation is $s=\frac{1}{2}$.
Our aim is to use probabilities to prove that this equation is almost surely (with regard to a certain measure) globally well-posed in subcritical spaces $H^{\sigma} \times H^{\sigma-1}$ with $\sigma \in[0,1 / 2)$.

## Result

Theorem There exist probability measures $\mu$ on spaces of low regularity such that $\mu\left(\mathrm{H}^{1 / 2} \times \mathrm{H}^{-1 / 2}\right)=0$ and for $\mu$-almost every $\left(f_{0}, f_{1}\right)$, the cubic wave equation with initial datum $\left(f_{0}, f_{1}\right)$ has a unique global solution in $L(t)\left(f_{0}, f_{1}\right)+C\left(\mathbb{R}, H^{1}\left(\mathbb{R}^{3}\right)\right)$ where $L(t)$ is the flow of the linear wave equation $\partial_{t}^{2}-\Delta=0$.

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## Compactification

The first step is to use the Penrose transform (conformal) to turn the problem on $\mathbb{R}^{3}$ into a problem on the sphere $S^{3}$ :

$$
\left\{\begin{array}{l}
\partial_{T}^{2} u+\left(1-\Delta_{S^{3}}\right) u+u^{3}=0 \\
u_{T T=0}=u_{0}, \partial_{T} u_{\mid T=0}=u_{1}
\end{array} .\right.
$$

Remark: This step is probably unnecessary. Though, skipping it implies using objects that seem less natural or at least not canonical.
The transform that maps $\left(u_{0}, u_{1}\right)$ to $\left(f_{0}, f_{1}\right)$ is an isometry between $H^{s} \times H^{s-1}$ of the sphere and $\mathcal{H}_{0}^{s} \times \mathcal{H}_{1}^{s-1}$ of $\mathbb{R}^{3}$ where $\mathcal{H}_{i}^{s}$ is very similar to $H^{s}$. In particular, if $\left(u_{0}, u_{1}\right)$ is not in $\mathrm{H}^{1 / 2} \times \mathrm{H}^{-1 / 2}$ then ( $f_{0}, f_{1}$ ) can not be in critical or super critical spaces.
Remark 2 : The existence of a solution of this compact equation gives the existence of a solution on $\mathbb{R}^{3}$. Uniqueness has to be treated separately.

## Reduction

The second step is to reduce the equation on $u$ on an equation on $v=u-U(T)\left(u_{0}, u_{1}\right)$ where $U(T)$ is the flow of the linear equation $\partial_{T}^{2}+1-\Delta_{S^{3}}=0$. We get

$$
\partial_{T}^{2} v+\left(1-\Delta_{S^{3}}\right) v+\left(U(T)\left(u_{0}, u_{1}\right)+v\right)^{3}=0
$$

with initial datum $v_{\mid T=0}=v_{0}=0$ and $\partial_{T} v_{\mid T=0}=v_{1}=0$.

## Local well-posedness

The Duhamel form of this equation is given by :
$v(T)=U(T)\left(v_{0}, v_{1}\right)-\int_{0}^{T} \frac{\sin ((T-\tau) \sqrt{1-\Delta})}{\sqrt{1-\Delta}}\left(U(\tau)\left(u_{0}, u_{1}\right)+v(\tau)\right)^{3} d \tau$.
The local theory yields that the Cauchy problem associated with this equation is well-posed in $H^{1}$ as soon as $v_{0} \in H^{1}, v_{1} \in L^{2}$ and $\frac{1}{\left(1+T^{2}\right)^{1 / 3}} U(T)\left(u_{0}, u_{1}\right) \in L_{T}^{3}, L^{6}\left(S^{3}\right)$.

## Global theory on $v$

We use energy estimates with

$$
\mathcal{E}(T)=\int_{S^{3}}\left(\partial_{T} v\right)^{2}+\int v(1-\Delta) v+\frac{1}{2} \int v^{4} .
$$

Gronwall lemma yields

$$
\mathcal{E}(T) \lesssim\left(\int_{0}^{T}\left\|U(\tau)\left(u_{0}, u_{1}\right)\right\|_{L^{6}}^{3} d \tau\right) e^{c \int_{0}^{T}\left(\left\|U(\tau)\left(u_{0}, u_{1}\right)\right\|_{L^{6}}^{2}+\left\|U(\tau)\left(u_{0}, u_{1}\right)\right\|_{L^{\infty}}\right) d \tau} .
$$

We have global well posedness in $U(T)\left(u_{0}, u_{1}\right)+C\left(\mathbb{R}, H^{1}\right)$ as soon as $U(T)\left(u_{0}, u_{1}\right)$ belongs to $L_{\text {loc }, T}^{1}, L^{\infty}\left(S^{3}\right)$.

## Conditions on the measure

We want to find a non trivial measure $\rho$ on the topological $\sigma$-algebra of $H^{\sigma} \times H^{\sigma-1}$ such that:

- $\rho\left(H^{1 / 2} \times H^{-1 / 2}\right)=0$,
- $\frac{1}{\left(1+T^{2}\right)^{1 / 3}} U(T)\left(u_{0}, u_{1}\right) \in L_{T}^{3}, L^{6}\left(S^{3}\right)$,
- $U(T)\left(u_{0}, u_{1}\right) \in L_{\mathrm{loc}, T}^{1}, L^{\infty}\left(S^{3}\right)$.

For this, we will randomize the initial data.

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## Notations

- $\left(e_{n, k}\right)_{n, k}$ is a $L^{2}$ orthogonal basis composed of spherical harmonics : - $\Delta_{S^{3}} e_{n, k}=n^{2} e_{n, k}, 1 \leq k \leq(n+1)^{2}$,
- $\left(a_{n, k}\right)_{n, k},\left(b_{n, k}\right)_{n, k}$ are two sequences of independent real Gaussian variables of law $\mathcal{N}(0,1)$ in a probability space $\Omega, \mathcal{A}, \mathbb{P}$,
- $\bar{u}_{0}=\sum \lambda_{n, k} e_{n, k}$ belongs to $H^{\sigma}$ for some $\sigma \in(0,1 / 2)$ but does not belong to $H^{1 / 2}$,
- $\bar{u}_{1}=\sum \mu_{n, k} e_{n, k}$ belongs to $H^{\sigma-1}$ (for the same $\sigma$ ) but not to $H^{-1 / 2}$.

Remark : The Gaussian condition can be released. We can take $\left(a_{n, k}\right)_{n, k},\left(b_{n, k}\right)_{n, k}$ two sequences of i.i.d random variables that satisfy : "there exists $c$ such that for all $\gamma$ and all $(n, k)$

$$
E\left(e^{\gamma a_{n, k}}\right), E\left(e^{\gamma b_{n, k}}\right) \leq e^{c \gamma^{2}}
$$

We can also take $\sigma=0$.

## Randomization

We then build two random variables :

$$
\begin{aligned}
& u_{0}(\omega)=\sum_{n, k} \lambda_{n, k} a_{n, k}(\omega) e_{n, k} \\
& u_{1}(\omega)=\sum_{n, k} \mu_{n, k} b_{n, k}(\omega) e_{n, k}
\end{aligned}
$$

We then have a measure $\rho$ on $H^{\sigma} \times H^{\sigma-1}$, the image measure of $\mathbb{P}$ by $\left(u_{0}, u_{1}\right)$, that is

$$
\rho\left(A_{0} \times A_{1}\right)=\mathbb{P}\left(u_{0}^{-1}\left(A_{0}\right)\right) \mathbb{P}\left(u_{1}^{-1}\left(A_{1}\right)\right) .
$$

Using the continuity of the space-time compactification on the initial datum, we can define the image measure $\mu$ of $\rho$ by this transform and get back on the Euclidean space this way.
For almost all $\omega \in \Omega, u_{0}(\omega)$ does not belong to $H^{1 / 2}$ and $u_{1}(\omega)$ does not belong to $H^{-1 / 2}$. For $\mu$ almost all $\left(f_{0}, f_{1}\right)$, $f_{0}$ does not belong to $H^{1 / 2}$ and $f_{1}$ does not belong to $H^{-1 / 2}$.

## Properties of the randomization

With a good choice of $e_{n, k}$ and $\sigma>0$, we have for all $(p, q) \in[1, \infty)$
$\rho\left(\left\{\left(\tilde{u}_{0}, \tilde{u}_{1}\right) \in H^{\sigma} \times H^{\sigma-1} \left\lvert\,\left\|\frac{1}{\left(1+T^{2}\right)^{1 / p}} U(T)\left(\tilde{u}_{0}, \tilde{u}_{1}\right)\right\|_{L_{T}^{p}, W^{\sigma, q}\left(S^{3}\right)}<\infty\right.\right\}\right)=1$.
In other words, for $\rho$ almost every $\left(\tilde{u}_{0}, \tilde{u}_{1}\right)$ taken in $H^{\sigma} \times H^{\sigma-1}$, the function $\frac{1}{\left(1+T^{2}\right)^{1 / p}} U(T)\left(\tilde{u}_{0}, \tilde{u}_{1}\right)$ belongs to $L_{T}^{p}, W^{\sigma, q}\left(S^{3}\right)$ and since $\sigma>0$ to $L_{T}^{p}, L^{\infty}\left(S^{3}\right)$.

Hence, $\rho$ almost every ( $\tilde{u}_{0}, \tilde{u}_{1}$ ) is a good candidate to be an initial datum of the equation on the sphere.

We prove that

$$
I:=\left\|\frac{1}{\left(1+T^{2}\right)^{1 / p}} D^{\sigma} U(T)\left(u_{0}, u_{1}\right)\right\|_{L_{\omega}^{r}, L_{T}^{p}, L q\left(S^{3}\right)}
$$

is finite with $r=\max (p, q)$ and $D=\left(1-\Delta_{S^{3}}\right)^{1 / 2}$.
The Minkowski inequality yields

$$
I \leq\left\|\frac{1}{\left(1+T^{2}\right)^{1 / p}} D^{\sigma} U(T)\left(u_{0}, u_{1}\right)\right\|_{L_{T}^{p}, L^{q}\left(S^{3}\right), L_{\omega}^{r}} .
$$

Then, at $x \in S^{3}$ and $T$ fixed, with $\langle n\rangle=\left(1+n^{2}\right)^{1 / 2}$

$$
\begin{aligned}
& \frac{1}{\left(1+T^{2}\right)^{1 / \rho}} D^{\sigma} U(T)\left(u_{0}, u_{1}\right)=\sum_{n, k} \frac{1}{\left(1+T^{2}\right)^{1 / \rho}}\langle n\rangle^{\sigma} \times \\
& \quad\left(\cos (\langle n\rangle T) a_{n, k} \lambda_{n, k} e_{n, k}(x)+\frac{\sin \langle\langle \rangle\rangle T)}{\langle n\rangle} b_{n, k} \mu_{n, k} e_{n, k}(x)\right)
\end{aligned}
$$

is a Gaussian variable as a linear combination of independent Gaussian variables. Hence, its $L_{\omega}^{r}$ norm is bounded by $C \sqrt{r}$ times its $L_{\omega}^{2}$ norm.

It gives

$$
\begin{aligned}
\left\|\frac{1}{\left(1+T^{2}\right)^{1 / p}} D^{\sigma} U(T)\left(u_{0}, u_{1}\right)\right\|_{L_{\omega}} \lesssim & \sqrt{r} \frac{1}{\left(1+T^{2}\right)^{1 / p}}\left(\sum _ { n , k } \left(\langle n\rangle^{2 \sigma}\left|\lambda_{n, k}\right|^{2}+\right.\right. \\
& \left.\left.\langle n\rangle^{2 \sigma-2}\left|\mu_{n, k}\right|^{2}\right)\left|e_{n, k}(x)\right|^{2}\right)^{1 / 2} .
\end{aligned}
$$

It remains to take the $L_{T}^{p}, L^{q}\left(S^{3}\right)$ of the right hand side of the inequality.

$$
I \leq C \sqrt{r}\left(\sum_{n, k}\left(\langle n\rangle^{2 \sigma}\left|\lambda_{n, k}\right|^{2}+\langle n\rangle^{2 \sigma-2}\left|\mu_{n, k}\right|^{2}\right)\left\|e_{n, k}\right\|_{L^{q}}^{2}\right)^{1 / 2} .
$$

We need a $L^{2}$ basis $\left(e_{n, k}\right)_{n, k}$ uniformly bounded in $L^{9}$.

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## A measure on basis

(Technique and result by Burq and Lebeau)
The idea is to build a measure on the basis of $L^{2}$ composed of spherical harmonics such that the probability that the basis is uniformly bounded in $L^{p}$ is non 0 .

We consider the $L^{2}$ orthonormal basis of spherical harmonics of degree $n$ as the orthogonal group $O_{N_{n}}(\mathbb{R})$ where $N_{n}$ is the dimension of spherical harmonics of degree $n: N_{n}=(n+1)^{2}$.

Take $v_{n}$ the Haar measure on $O_{N_{n}}(\mathbb{R})$ and

$$
v=\otimes_{n \in \mathbb{N}} v_{n}
$$

To evaluate probabilities on one element $b_{n, k}$ of the basis, we have to take the $k$-th column of the matrix associated to $\left(b_{n, j}\right)_{j}$.

The image measure of $v_{n}$ by the map that takes the $k$-th column is (thanks to invariances) the uniform probability measure on the sphere $S^{N_{n}}: p_{N_{n}}$.

## Measure concentration phenomenon

The Lipschitz-continuous functions $F$ from $S^{N}$ to $\mathbb{R}$ concentrate on their median $M_{F}$ in the sense that

$$
p_{N}\left(\left\{x| | F(x)-M_{F} \mid \geq R\right\}\right) \leq 2 e^{-(N-1) R^{2} /\left(2\|F\|_{\text {lip }}^{2}\right)}
$$

with

$$
|F(x)-F(y)| \leq\|F\|_{l i p}\|x-y\|_{2} .
$$

The $L^{p}$ norm is Lipschitz continuous on spherical harmonics of degree $n$ :

$$
\|x-y\|_{L^{p}} \leq C n^{1-1 / p}\|x-y\|_{L^{2}}
$$

and its median $M_{p}$ is bounded by $C \sqrt{p}$ (independent from $n$ ). Both are consequences of the fact that $S^{3}$ has a finite volume and that for all $x, y$ in $S^{3}$, there exists a transformation $R$ on $S^{3}$ that preserves the metrics such that $x=R y$.

We get then

$$
\begin{aligned}
v_{n}\left(\left\|b_{n, k}\right\|_{L^{p}}-M_{p} \geq R\right) & \leq p_{n}\left(\left|\|x\|_{L^{p}}-M_{p}\right| \geq R\right) \\
& \lesssim e^{-c R^{2} n^{4 / p}} .
\end{aligned}
$$

By summing over $k$

$$
v_{n}\left(\exists k \mid\left\|b_{n, k}\right\|_{L^{p}}-M_{p} \geq R\right) \lesssim n^{2} e^{-c R^{2} n^{4 / p}}
$$

and over $n$

$$
v\left(\exists n, k \mid\left\|b_{n, k}\right\|_{L^{p}}-M_{p} \geq R\right) \leq C_{p} R^{-2} .
$$

## Conclusion

We have used probabilities in two ways:

- randomizing the initial datum makes it almost surely in $L^{p}$ spaces,
- randomizing the basis enables us to take one uniformly bounded in $L^{p}$.
If $\sigma$ (the regularity of the ID) is 0 , then we do not have $U(T)\left(u_{0}, u_{1}\right)$ almost surely in $L^{\infty}$ but being careful with the choice of the basis and using a bootstrap argument in the energy estimates, we still have the same result as before.

